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INVESTIGATION OF AN ALL-MOVABLE CONTROL SURFACE AT

A MACH NUMBER OF 6.86 FOR POSSIBLE FLUTTER

By William T. Lauten, Jr., Gilbert M. Levey, and William O. Armstrong

Langley Aeronautical Laboratory
Langley Field, Va.

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RESEARCH MEMORANDUM

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SUMMARY

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Results of tests for possible flutter of a dynamically and elastically scaled model of a proposed all-movable horizontal tail surface for the North American X-15 airplane are presented herein. Tests at a Mach number of 6.86 were made on the scaled model and on several other configurations having lower stiffnesses. No flutter was obtained. Flexibility influence coefficients and calculated modes and frequencies for the weakest configuration are presented. Calculations of flutter speed of the weakest configuration (piston-theory aerodynamic forces, calculated mode shapes, and experimentally determined frequencies in a modal-analysis calculation scheme being used) yielded a flutter speed approximately four times as high as the velocity obtained in the tests.

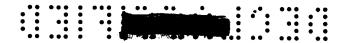
INTRODUCTION

Recent developments in aircraft and missiles with all-movable controls have led to increased interest concerning the flutter of such plan forms. Current design trends indicate the need for flutter information on all-movable plan forms at hypersonic speeds. Reference 1 reports tests of a research program in which a rectangular-plan-form, all-movable control was tested at a Mach number of 6.86.

As a part of a program of flutter testing of the various surfaces of the North American X-15 airplane, tests were made in the Langley ll-inch hypersonic tunnel on a 1/12-size model, dynamically and elastically scaled on the basis of dynamic pressure, of a proposed horizontal tail surface. In addition to the scaled model, several configurations with reduced panel and mounting stiffnesses were tested.

Presented herein are the test conditions during the tunnel runs, the structural characteristics of the various configurations, and, for the





weakest configuration, the measured flexibility influence coefficients and mode shapes and frequencies calculated from the influence coefficients. The results of flutter calculations on the weakest configuration are included. These calculations were made by using piston-theory aerodynamic forces (ref. 2), the first three calculated mode shapes, and the corresponding experimentally determined natural frequencies.

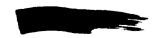
SYMBOLS

а	velocity of sound, fps
f	frequency, cps
g	damping coefficient, theoretical value needed to produce flutter
M	Mach number
q	dynamic pressure, lb/sq ft
ρ	air density, slugs/cu ft
Subscripts:	
1,2,3,4	indicate natural vibration mode in order of ascending frequency
exp	experimental
calc	calculated

MODEL DESCRIPTION

The model of the all-movable control surface was 1/12 scale with an exposed-surface aspect ratio of approximately 2.5, a taper ratio of 0.305, and a sweep angle of 45° at the quarter-chord line. The airfoil was an NACA 66AOO5 modified so that it was 1 percent thick at the trailing edge with a straight-line fairing to the point of tangency. The airfoil ordinates are listed in table I.

A top-view drawing of the model mounted in the test section is shown in figure 1(a). Figure 1(b) shows a three-dimensional sketch of the detail of the spindle and spring restraints. The model was supported in its base block by means of two flexure springs attached to the spindle in



such a way that the flexure springs provided the same restraint points as the spindle bearings in the prototype. The flexure spring pivot was at 25 percent of the mean aerodynamic chord. A third spring, similar to the flexure springs, was mounted ahead of the model spindle and attached to the spindle with a screw as shown in figure 1(b). This spring provided additional pitch stiffness and is referred to as the pitch spring. The stiffness of the three springs combined to simulate the actuator stiffness of the prototype.

Several configurations were tested. The basic model was dynamically and elastically scaled from the prototype, both in panel stiffness and spindle restraint, on the basis of dynamic pressure. The model wing consisted of a solid aluminum spar, on which the thickness was varied to obtain the desired stiffness distribution, and five hollow aluminum streamwise segments, as shown in figure 2. The spar had an integral spindle for mounting. Each segment was fitted to the spar and fastened in place with two screws in such a way that the wing stiffness was determined by the spar stiffness with the segments contributing a negligible amount. This model is referred to as model I in table II.

Several weaker configurations were tested. Model II had the same construction as model I except that a hollow spar was used which had about two-thirds the stiffness of the spar of model I. Model III was the same as model II except the slots between the segments were covered over with fiber-glass tape. This resulted in a somewhat stiffer panel. Model IV was identical to model I except that the spar was drilled out with 0.1875-inch-diameter holes 0.244 inch on center in order to reduce the stiffness further. This method of controlling stiffness is discussed in detail in reference 3. Model V was a modification of model IV in that the holes were enlarged to 0.204 inch and the spindle was milled so that its calculated stiffness was cut in half. Figure 3 shows the final spar and spindle for model V. The small holes in the edges and tip were an attempt to obtain a more uniform stiffness and mass distribution. The frequency spectrum of model V indicated that it had about one-fourth the stiffness of model I.

The model base block, shown in figure 1, served as a model mount and also as a spacer to support the model in the airstream beyond the tunnel-wall boundary layer. A reflection plane, also shown in figure 1, was attached to the model base block just inboard of the wing root. A photograph of the model mounted in the tunnel is shown in figure 4.





INSTRUMENTATION

A recording oscillograph was used to obtain continuous records during each test from strain gages oriented on the model spar to record strain about two axes (primarily panel bending and torsion) and from a strain gage mounted on the pitch spring. Simultaneously recorded were the outputs from a thermocouple and a pressure cell from which tunnel stagnation temperature and pressure could be determined. Motion pictures at a speed of 128 frames per second were taken of the model during each tunnel run.

LABORATORY MEASUREMENTS

As a preliminary to each tunnel run the frequencies and node lines for the first four natural vibration modes were determined by use of an acoustic shaker. Figure 5 shows a typical set of node lines.

In addition to the natural frequencies of all models, flexibility influence coefficients were determined prior to testing for models IV and V (runs 9 and 11). The location of the points at which the models were loaded is shown in figure 6. The influence coefficients for model V (the weakest model) are tabulated in table III with the calculated mode shapes and frequencies for the first three natural modes. Also listed are the experimentally determined frequencies.

The wing deflections were measured electrically by an array of differential transformers. These transformers were connected to a null balance indicator with a visual indication which could be read to 0.0001 inch. This system was accurate to ±0.001 inch from 25 to 75 percent of full range; however, for reading small differences, for example, the difference between 0.0550 and 0.0560 inch, the readings could be repeated to at least ±0.0001 inch.

The mass, center of gravity, and moment of inertia of each of the five wing segments were determined experimentally. The corresponding quantities for the portion of the spar of model V associated with each segment were calculated. Then, in order to meet the requirement for correct total mass, center of gravity, and moment of inertia of any streamwise segment while still associating two masses with the loading points of that segment, a fictitious mass was assigned to each segment at the locations shown in figure 6. The displacement associated with each such mass was found by linear interpolation between the loading points of that segment. This results in dynamic coupling terms in the mass matrix. (See table IV.) These terms are similar to those introduced by Rodden in reference 4 and, while satisfying the three known conditions



of mass, center of gravity, and moment of inertia of a strip, give the proper kinetic energy of the vibrating strip.

TEST PROCEDURE

The tests were made in the Langley 11-inch hypersonic tunnel by using a single-step, two-dimensional Invar nozzle. With this nozzle, the tunnel operates at an average Mach number of 6.86. A description of the tunnel is given in reference 5 and a preliminary calibration of the Invar nozzle is included in reference 6.

In an effort to determine a flutter boundary, the density in the test section was increased during each run by gradually increasing the stagnation pressure from 5 atmospheres to a peak value of about 38 atmospheres. For normal runs, about 30 seconds were required for the dynamic pressure to reach a maximum value.

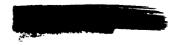
In order to alleviate the danger of damage to the model by the starting and stopping transients, restraining pins, operated by a lever outside the tunnel, were inserted into the root section until the starting shock passed through the test section, were retracted during the increase to maximum pressure, and were reinserted before the tunnel closed down.

There was deviation from the normal testing procedure on three of the tests. In run 6, the tunnel was started at maximum stagnation pressure - 42 atmospheres; in run 7, the wing was preloaded at the tip and released suddenly when the tunnel reached near maximum pressure; and in run 10, the model was tested at an angle of attack of about $2\frac{1}{2}$. None of these variations had any apparent effect on the model stability.

RESULTS AND DISCUSSION

A summary of the test results is given in table II, which lists the first four natural frequencies, test-section density, speed of sound, and dynamic pressure at maximum pressure test condition and presents brief explanatory remarks for each run. No flutter was obtained on any of the configurations tested.

Since reference 1 reported good agreement between calculated and experimental flutter speed at the test Mach number, it was felt worthwhile to make the same calculation on model V reported here. The calculations were made by using the aerodynamic forces derived from piston



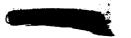


theory (ref. 2), the first three coupled vibration modes as calculated from the influence coefficients, and the experimentally determined frequencies for these modes in a modal-analysis calculation scheme. effects of thickness were included. This calculation yielded a result that may be interpreted alternatively as follows: (1) The flutter speed was about four times the maximum speed reached in the test or (2) the stiffness of model V would have to be decreased by a factor of approximately 16 before flutter could be expected to occur. This reduction would have been impractical with the model as originally constructed. The results of these calculations are presented in figure 7 in the form of a plot of velocity against damping coefficient for the only mode of vibration that had an instability. The damping coefficient g was assumed to have the same value for all modes and provides for theoretical or actual damping forces as illustrated, for example, in section 8.1 of reference 7.

CONCLUDING REMARKS

Flutter tests at a Mach number of 6.86 of a dynamically and elastically scaled model of a sweptback, all-movable horizontal tail proposed for the North American X-15 airplane and on several configurations having lower stiffnesses are reported. The spectrum of natural vibration frequencies indicates that the weakest configuration was one-fourth as stiff as the stiffest configuration. No flutter was obtained.

Langley Aeronautical Laboratory, National Advisory Committee for Aeronautics, Langley Field, Va., February 12, 1958.



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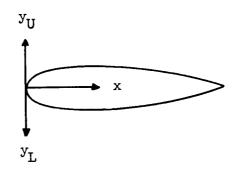




TABLE I ORDINATES FOR NACA 66A005 (MODIFIED) AIRFOIL

Ordinates in percent of chord

х	$y_U = y_L$
0 .10 .25 .50 .75 1.25 2.50 5.00 7.50 10.00 25.00 25.00 30.00 40.00 45.00 50.00 60.00 65.00 a67.00	0 .408 .531 .590 .650 .791 1.048 1.270 1.460 1.766 2.001 2.182 2.318 2.416 2.476 2.476 2.488 2.476 2.488 2.476 2.176 2.176 2.085 .500





^aStraight-line fairing from 67 to 100 percent chord.



EXPERIMENTAL DATA

Run	Model		. – i	fz,		ρ,	а,	q,		Remarks
		cps	cps	cps	cps	slugs/cu ft	ıps	TD/sq	Ιτ	
1	I	88	263	293	461	0.000128	506	769		Solid spar and stiffest pitch spring
2	IA	88	258	289	475	.000110	528	724		Solid spar and weak- ened pitch spring
3	IB	87	254	287	479	.000111	521	7 07		Solid spar and no pitch spring
14	II	82	224	252	392	.000108	531	719		Hollow spar and no pitch spring for rest of runs
5	IIA	79	207	251	3 85	.000115	522	735		Hollow spar, reduced flexure stiffness
6	AII	7 9	207	251	3 85	.000134	500	793		Same as run 5, except run started at maximum density
7	TIA	79	207	251	385	.000129	512	794		Same as run 5, but model excited during run
8	III	83	227	309	580	.000128	480	692	!	Same as run 5, but slots taped over
9	IV	54	138	220	31 9	.000126	496	730		Solid spar drilled with 0.1875-inch- diameter holes
10	IV	54	138	220	3 19	.000119	500	702		Same as run 9 but model at angle of attack of 210
11	V	44	115	148	172	.000112	521	713		Holes in spar enlarged to 0.204 inch in diameter, and spindle stiffness decreased

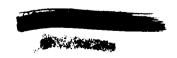


TABLE III

MEASURED FLEXIBILITY INFLIENCE COEFFICIENTS AND CALCULATED-MODE-SHAFE

VALUES AT STATIONS INDICATED IN FIGURE 6 FOR MODEL ${f v}$

0.0	0.0096 .0050 .0012 .0012 .0045 .0090 .0090 .0070
6	0.0145 .0093 .0097 .0097 .0097 .0056 .0166 .0166
8	0.0350 .01855 .0036 .0058 0074 0045 0045 0045 0045 0020
7	0.0537 .0290 .0091 .0019 .0175 .0175 .075 .0315 .0315
9	0.0886 0.0886 0.0077 0.0011 0.0111 0.0111 0.0296 0.0296
5	-0.0085 -0.0047 -0.0011 -0.0022 -0.0071 -0.016 -0.0074 -0.0074
4	-0.0018 0003 0002 0031 0034 0019 0010
8	420.0 420.0 600.0 600.0 600.0 600.0 600.0 600.0 600.0 600.0
0	0.0353 0.0250 0.0201 0.0017 0.0407 0.0200 0.0290 0.0290
Loading	100000000 1000000000000000000000000000
	9 7 8 9

Third mode	1.0000 .99770 .56289 .26615 .01802 52444 18950 .04130	163.4 148	
Second	1.0000 . 4.2942 . 3.2034 . 3.2034 . 7.2664 . 7.7199 . 98218 . 5.7259 . 4.7219 . 4.7219	110.0 115	
First	0.73297 .39308 .12749 03655 16073 1.0000 .68561 .43781 .24328	140.5 44	
Loading	10 00 10 00 pg	fcalc, cps fexp, cps	

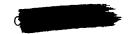
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TABLE IV

MASS MATRIX WITH DYNAMIC COUPLING TERMS

Mass in slugs

10	0	0	0	0	-0.0032467	0	0	0	0	0.0043711
6	0	0	0	-0.0021861	0	0	0	0	0.0032115	0
8	0	0	-0.0015243	0	0	0	0	0.0023088	0	0
7	0	-0.0010145	0	0	0	0	0.0016016	0	0	0
9	-0.0006818	0	0	0	0	0.0010095	0	0	0	0
5	0	0	0	0	0.0025899	0	0	0	0	-0.0032467
†	0	0	0	0.0019498	0	0	0	0	-0.0021861	0
3	0	0	0.0014018	0	0	0	0	-0.0015243	0	0
2	٥.	6016000.0	0	0	0	0	-0.0010145	0	0	0
ч	0.0005668	0	0	0	0	-0.0006818	0	0	0	0
Station	1	ณ	3	ⅎ	ζ.	9		ω	6	91



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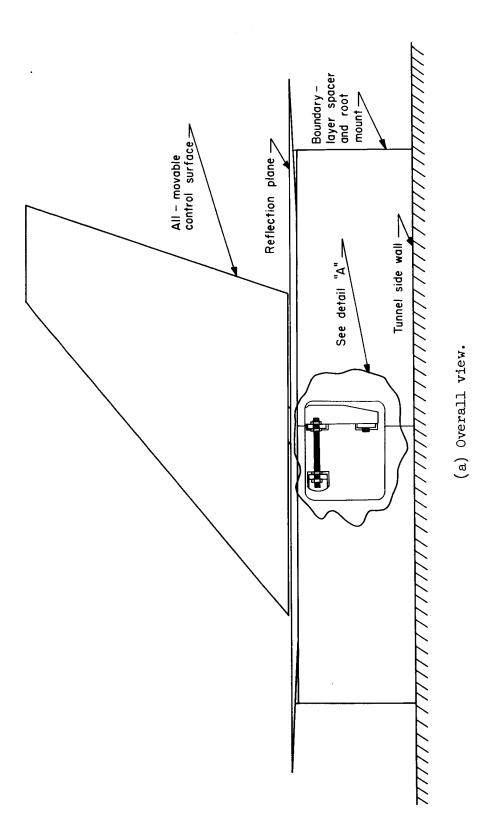
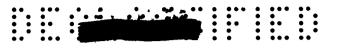
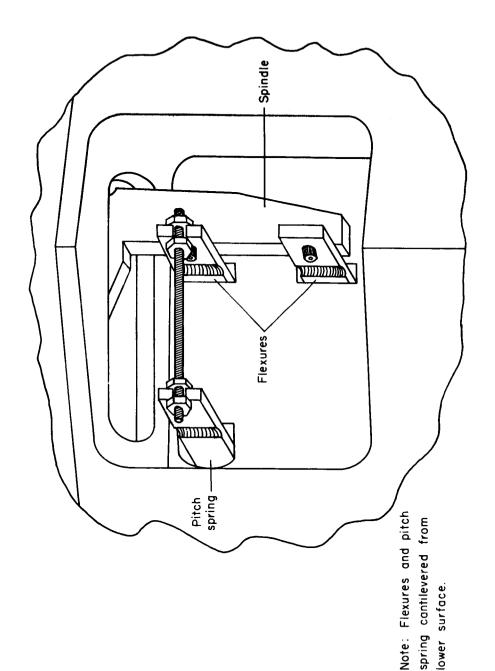


Figure 1.- Top view of wing as mounted in the tunnel.



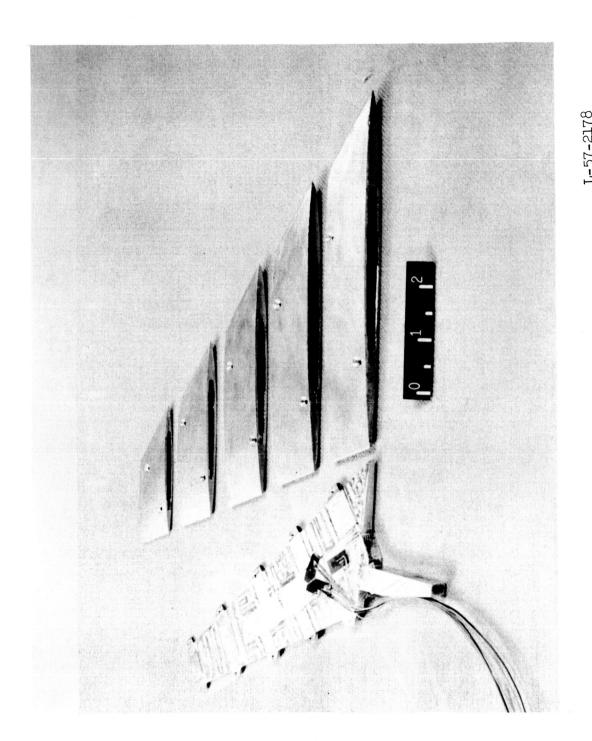


Detail "A"

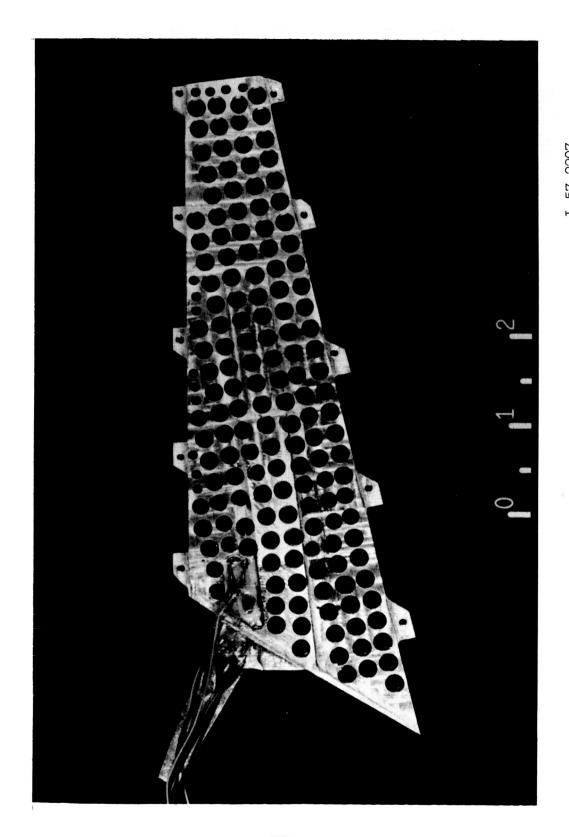
(b) Sketch showing detail of spindle support system.

Figure 1.- Concluded.

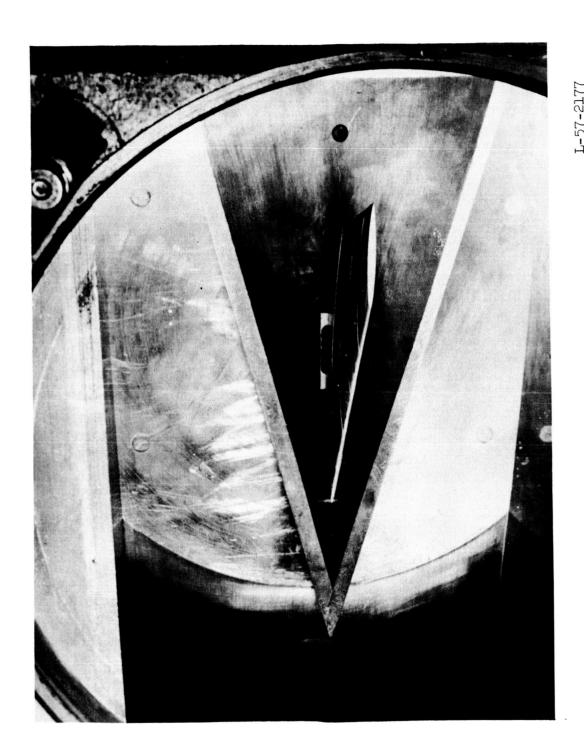




L-57-2178 Figure 2.- Photograph of solid-spar model with segments detached.

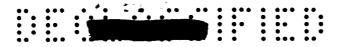


L-57-2207 Figure 3.- Photograph of solid spar showing hole pattern for model V_{\bullet}



L-57-2177 Figure 4.- Photograph of model mounted in tunnel showing reflection plane.





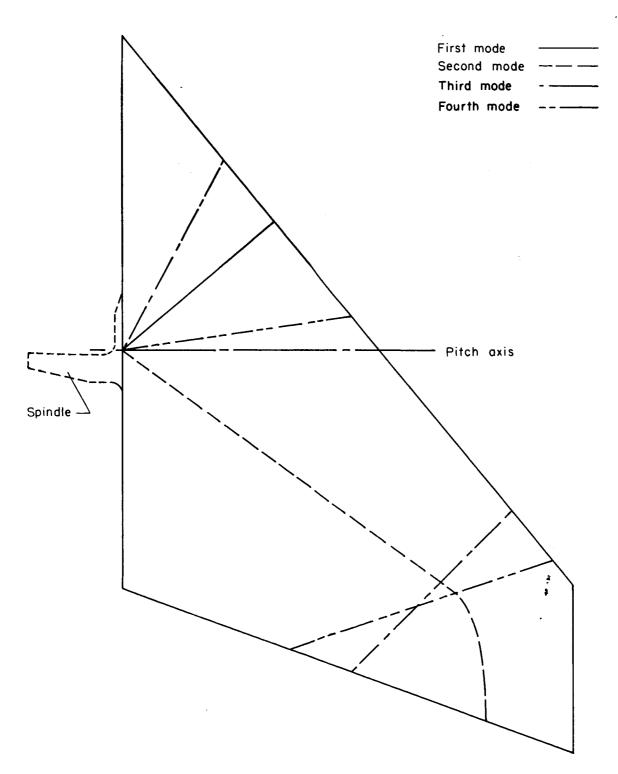


Figure 5.- Sketch of model showing typical nodal patterns.



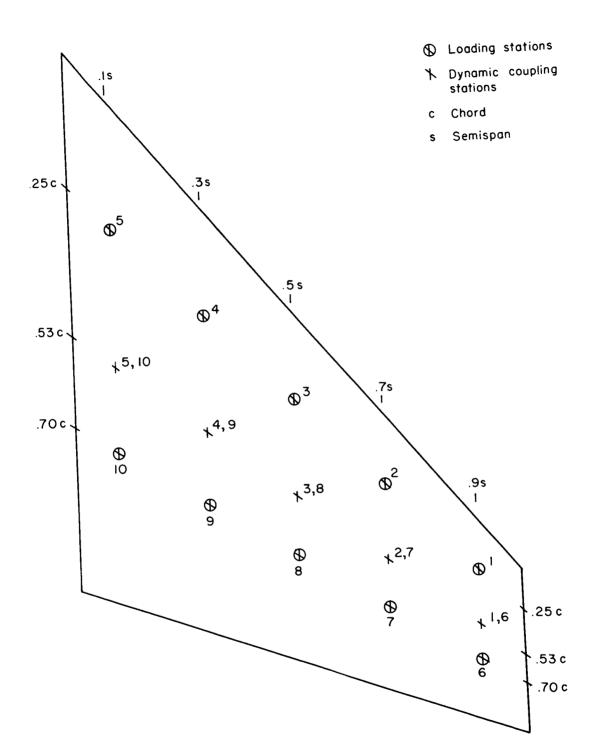
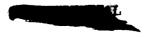
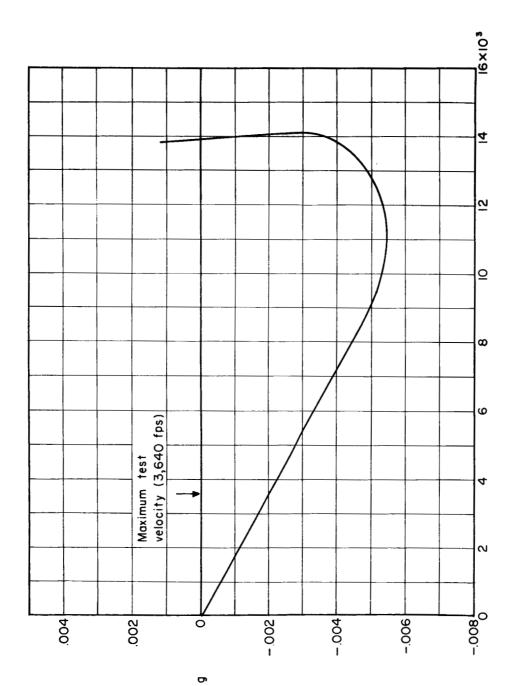


Figure 6.- Sketch of model showing points of load application and measurement of flexibility influence coefficients and the location of fictitious masses that are used in the analytical solution.





Velocity, fps

Figure 7.- Plot of velocity against damping coefficient for the only unstable vibration mode.